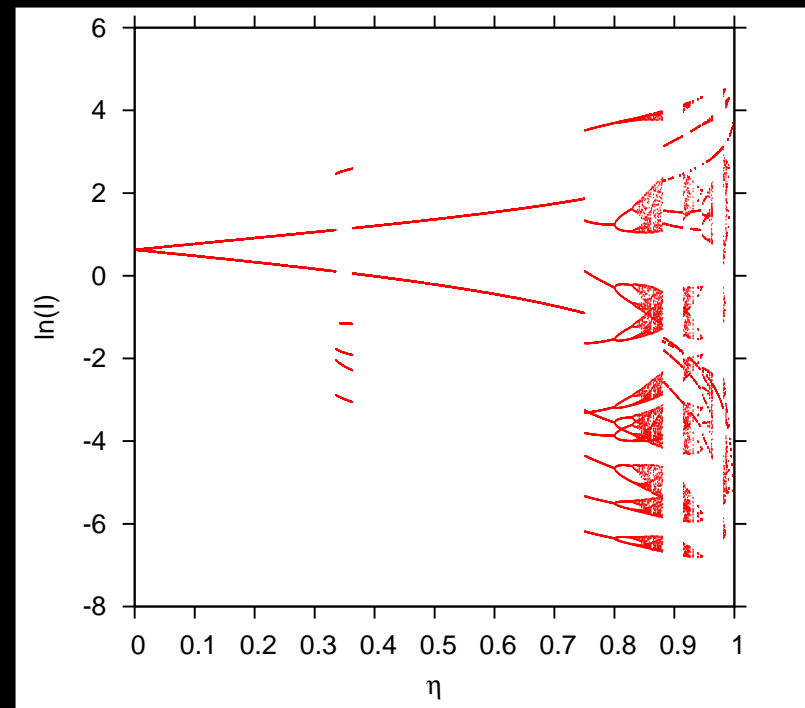
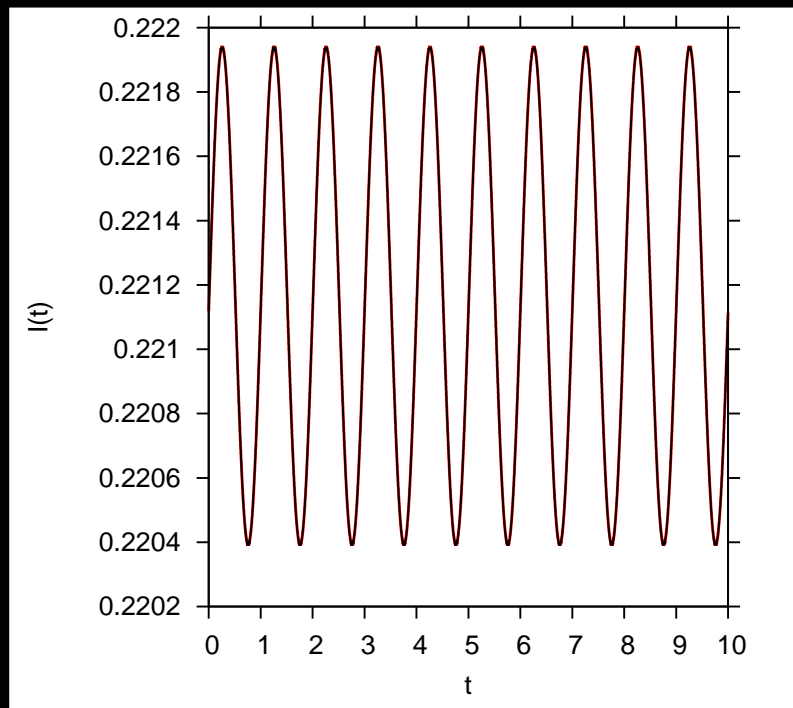


# The role of seasonality in vector-borne disease dynamics



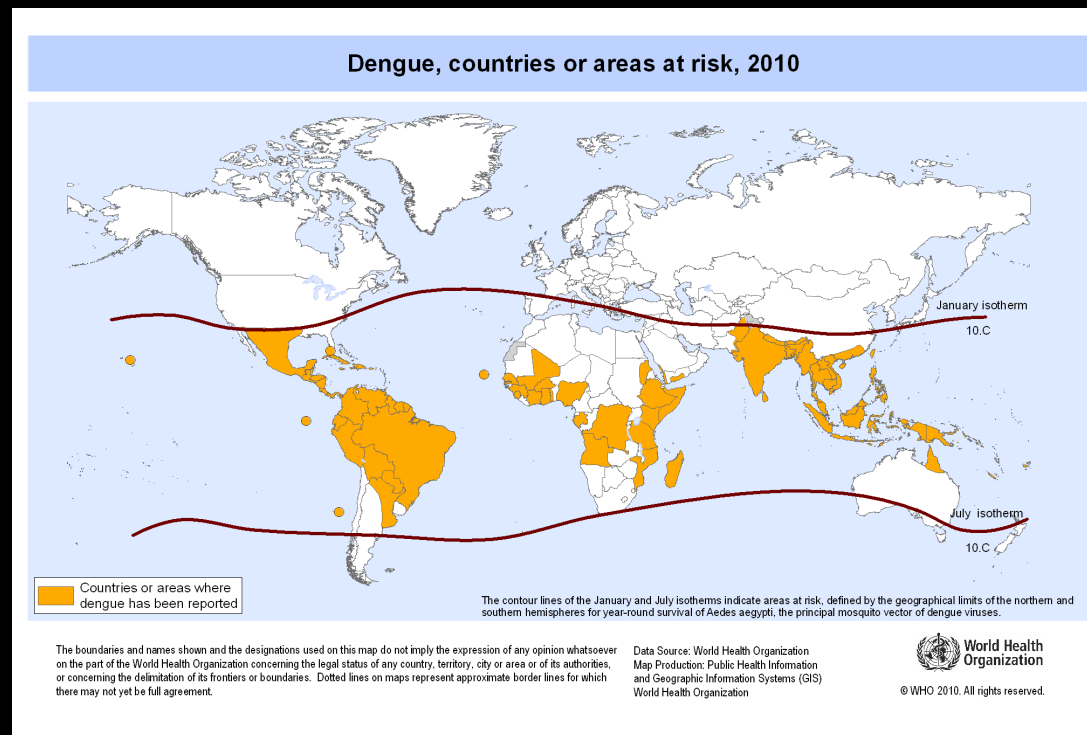
*Filipe Rocha, Maíra Aguiar and Nico Stollenwerk*

# *Dengue Fever Epidemiology*

Dengue is a mosquito-borne infection caused by an arbovirus with 4 serotypes DENV1-4.

The distribution is in tropical and subtropical areas. However, the disease is spreading to northern sites, being in the "gates of Europe".

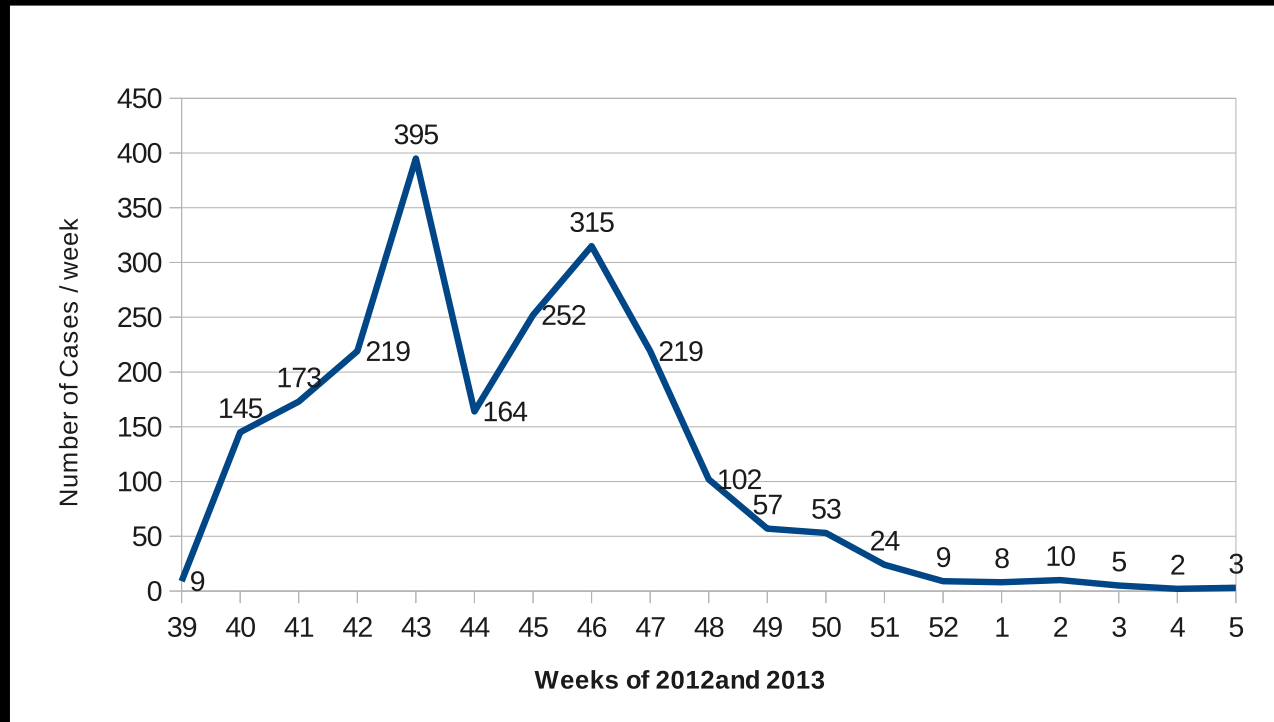
Recently the disease arrived to Madeira island and there is an outbreak with more than 2000 cases.



Worldwide dengue distribution in 2010 and areas at risk - Source: WHO (2012)

## *Dengue in Madeira*

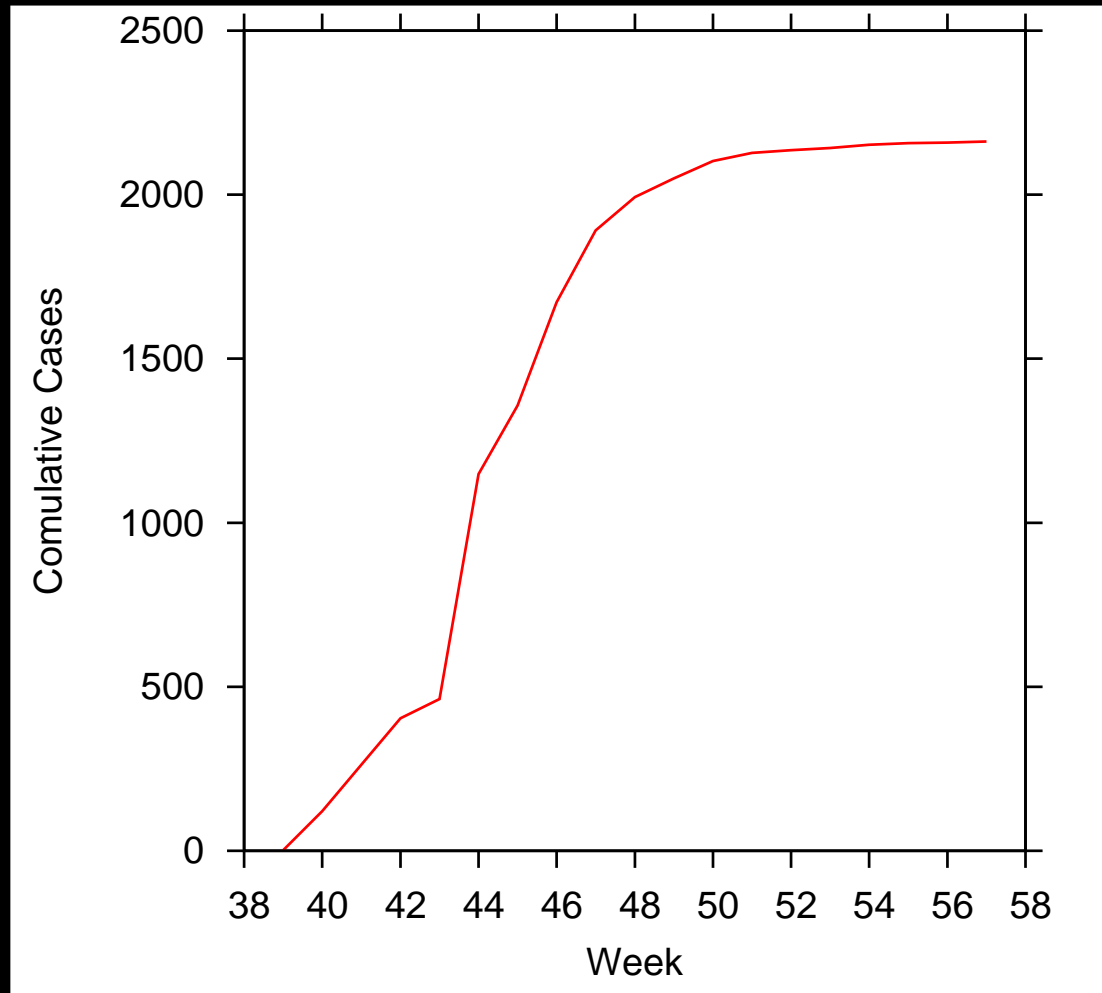
The outbreak started in early Autumn season and has been developing.



The virus seem to be the DENV-1, and people think that the disease was imported from the Americas (Brazil or Venezuela).

## *Dengue in Madeira*

Comulatively it has been reported 2164 cases.



# *Dengue in Madeira*

The outbreak is mainly in Funchal, but the disease is spreading through the island and surrounding islands.



Moreover, 78 cases of infected people were exported from the archipelago. Mainly in people from Portugal, but also from other countries such as UK, Germany, Sweden, France and Finland.

## *The vector*

The main vector is the *Aedes aegypti*, original from Africa, is now more distributed in Americas.

This species has been identified in Madeira island since 2005.

Other vector species is the *Aedes albopictus*, which is more distributed in Asia, Northern Africa and Europe.

This species has been identified to Spain, France, Italy, Croatia, Greece, between others.

Usually is verified an increase in number of mosquitos during the warmer seasons, specially in temperate regions.



## *Time-scale separation in SISUV*

The simplified version of the SISUV model, considering constant population size for human and mosquitos, is

$$\frac{d}{dt}I = \frac{\beta}{M}(N - I)V - \alpha I$$

$$\frac{d}{dt}V = \frac{\vartheta}{N}(M - V)I - \nu V$$

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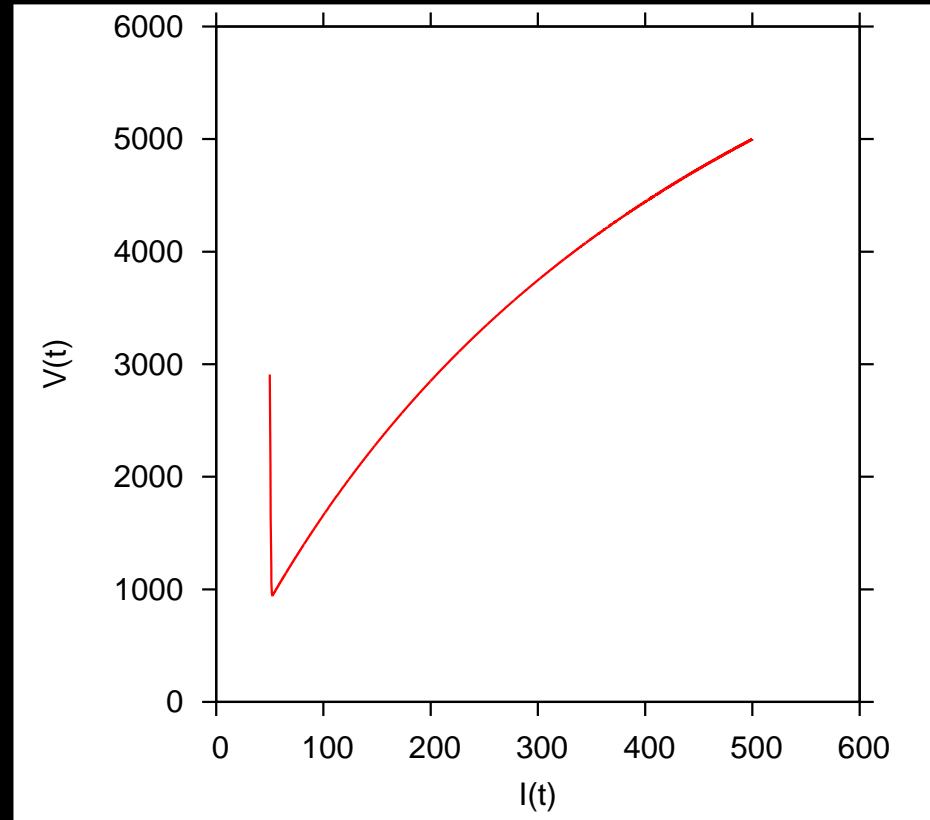
$$\begin{aligned}\frac{d}{dt}I &= \frac{\beta}{M}(N - I)V - \alpha I \\ \frac{d}{dt}V &= \frac{\vartheta}{N}(M - V)I - \nu V\end{aligned}$$

However it is noticed that the mosquitos' dynamics is faster than the humans', so we modified the variables os mosquitos dynamics ( $\vartheta =: \frac{\bar{\vartheta}}{\varepsilon}$  and  $\nu =: \frac{\bar{\nu}}{\varepsilon}$ ) in order to put them in the same range of the human's.

$$\begin{aligned}\frac{d}{dt}I &= \frac{\beta}{M}(N - I)V - \alpha I \\ \frac{d}{dt}V &= \frac{1}{\varepsilon} \left( \frac{\bar{\vartheta}}{N}(M - V)I - \bar{\nu}V \right)\end{aligned}$$



## *Time-scale separation in SISUV*



Using the following parameters set:

$$\alpha = \frac{1}{10} \text{y}^{-1}, \beta = 2 \cdot \alpha, \nu = \frac{1}{10} \text{d}^{-1} = \frac{365}{10} \text{y}^{-1} \text{ and } \vartheta = 2 \cdot \nu .$$

## *Time-scale separation in SISUV*

Considering the normal time scale given by  $t$  and the fast time scale given by  $\tau := \frac{t}{\varepsilon}$ , the general solution for the ODE system is:

$$\begin{aligned} I &= I_0 + \varepsilon I_1 + \varepsilon^2 I_2 + \mathcal{O}(\varepsilon^3) \\ V &= V_0 + \varepsilon V_1 + \varepsilon^2 V_2 + \mathcal{O}(\varepsilon^3) \end{aligned}$$

And for the slow time scale we obtain from the right hand side of the ODE system

$$\begin{aligned} \frac{dI}{dt} &= \varepsilon^0 \left( \frac{\beta}{M} (NV_0 - I_0V_0) - \alpha I_0 \right) + \varepsilon^1 \left( \frac{\beta}{M} (NV_1 - I_1V_0 - I_0V_1) - \alpha I_1 \right) + \mathcal{O}(\varepsilon^2) \\ \frac{dV}{dt} &= \frac{1}{\varepsilon} \left( \frac{\bar{\vartheta}}{N} (MI_0 - V_0I_0) - \bar{\nu}V_0 \right) + \varepsilon^0 \left( \frac{\bar{\vartheta}}{N} (MI_1 + V_1I_0 + V_0I_1) - \bar{\nu}V_1 \right) + \mathcal{O}(\varepsilon^1) \end{aligned}$$

## *Time-scale separation in SISUV*

Being  $\frac{dI}{d\tau} = \varepsilon \frac{dI}{dt}$  and  $\frac{dV}{d\tau} = \varepsilon \frac{dV}{dt}$ , if we substitute on the right hand of the ODEs we obtain

$$\frac{dI}{d\tau} = \varepsilon \underbrace{\left( \frac{\beta}{M} (N - I_0) V_0 - \alpha I_0 \right)}_{=\frac{dI_0}{d\tau}} + \mathcal{O}(\varepsilon^2)$$

$$\frac{dV}{d\tau} = \varepsilon^0 \underbrace{\left( \frac{\bar{\vartheta}}{N} (M - V_0) I_0 - \bar{\nu} V_0 \right)}_{=\frac{dV_0}{d\tau}} + \mathcal{O}(\varepsilon^1)$$

Or, for exactly  $\varepsilon = 0$ , the derivatives are:

$$\begin{aligned} \frac{dI_0}{d\tau} &= 0 \\ \frac{dV_0}{d\tau} &= \left( \frac{\bar{\vartheta}}{N} (M - V_0) I_0 - \bar{\nu} V_0 \right) \end{aligned}$$

## *Time-scale separation in SISUV*

As the infected has not fast time-scale, so  $\frac{dI_0}{d\tau} = 0$  and all values of  $I_0(\tau) = I_0(\tau_0)$ . So, substituting  $I_0(\tau)$  in  $\frac{dV_0}{d\tau}$  it is obtained:

$$\frac{dV_0}{d\tau} = - \left( \frac{\bar{\vartheta}}{N} I_0(\tau_0) + \bar{\nu} \right) V_0 + \frac{\bar{\vartheta}}{N} M I_0(\tau_0)$$

Which approaches very rapidly in an exponential way to its local stationary state:

$$V_0^* = \frac{\frac{\bar{\vartheta}}{N} I_0(\tau_0)}{\frac{\bar{\vartheta}}{N} I_0(\tau_0) + \bar{\nu}} \cdot M$$

## *Time-scale separation in SISUV*

Now to the slow dynamics:

$$\begin{aligned}\frac{dI_0}{dt} &= \left( \frac{\beta}{M} (NV_0 - I_0V_0) - \alpha I_0 \right) \\ \varepsilon \frac{dV_0}{dt} &= \left( \frac{\bar{\vartheta}}{N} (MI_0 - V_0I_0) - \nu V_0 \right)\end{aligned}$$

If we set  $\varepsilon = 0$ , we can obtain the equation of  $V_0(t)$ , for any time  $t$ :

$$V_0(t) = \frac{\frac{\bar{\vartheta}}{N} I_0(t)}{\frac{\bar{\vartheta}}{N} I_0(t) + \bar{\nu}} \cdot M$$

And now, finally, we can find the global stationary state:

$$I^* = \frac{\beta - \alpha \cdot \frac{\nu}{\bar{\vartheta}}}{\beta + \alpha} N \quad \text{and} \quad V^* = \frac{\beta - \alpha \cdot \frac{\nu}{\bar{\vartheta}}}{\beta \left(1 + \frac{\nu}{\bar{\vartheta}}\right)} M$$

## *Time-scale separation in SISUV*

The Jacobian matrix of the model is given by:

$$A = \begin{pmatrix} -\frac{\beta}{M} \cdot V^* - \alpha & \frac{\beta}{M} \cdot (N - I^*) \\ \frac{\bar{\vartheta}}{\varepsilon N} \cdot (M - V^*) & \frac{1}{\varepsilon} \left( -\frac{\bar{\vartheta}}{N} \cdot I^* - \bar{\nu} \right) \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$$

The eigenvalues of are given by:

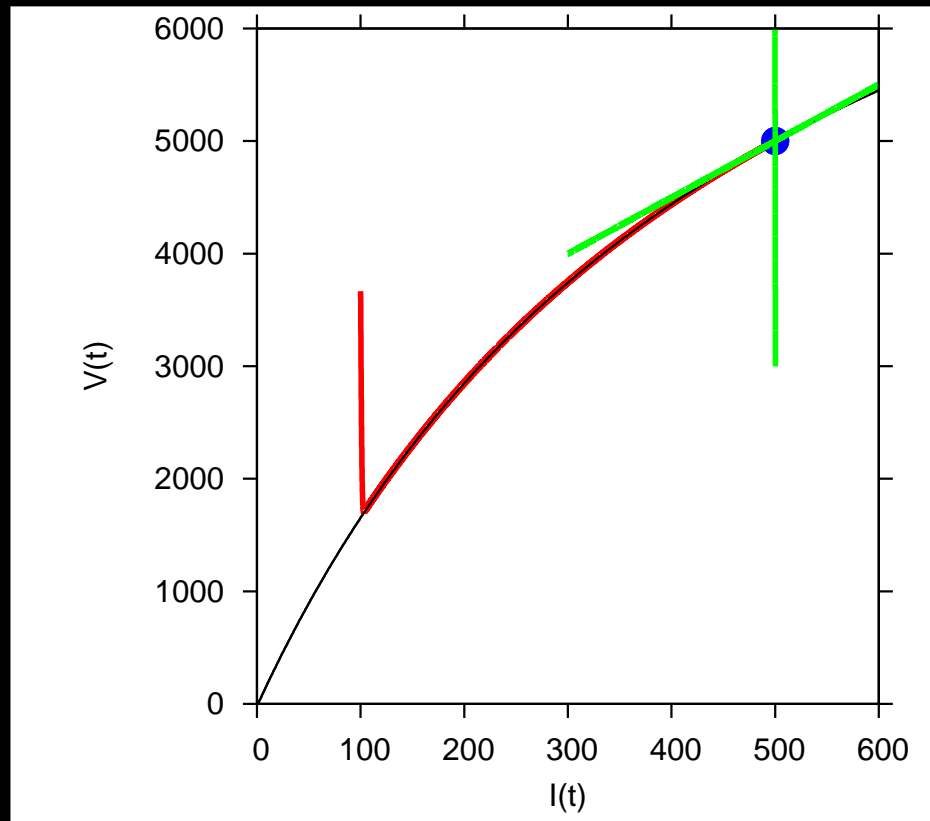
$$\lambda_{1/2} = \frac{(a + d)}{2} \pm \sqrt{\left(\frac{a + d}{2}\right)^2 - (ad - bc)}$$

And the numerical simulations shows that one is close to 0 and the other is large negative (  $\lambda_1 = 0$  and  $\lambda_2 = -73$ ).

And the general formula of eigenvectors is:

$$\underline{u}_i = \frac{1}{\sqrt{1 + \left(\frac{c}{d - \lambda_i}\right)^2}} \begin{pmatrix} 1 \\ -\frac{c}{d - \lambda_i} \end{pmatrix}$$

# *Time-scale separation in SISUV*



## *Center manifold analysis in SISUV*

Start by shifting the system  $(I, V)$  into a  $(z, w)$  system with the endemic fixed point at the origin:

$$\begin{aligned}z &:= I - I^* \\w &:= V - V^*\end{aligned}$$

Rearranging the system and considering the non-trivial stationary state as the origin of a  $(x, y)$  system and the eigendirections as coordinate axis. This transformation is done considering:

$$\underline{x} := T^{-1} \underline{z}$$

Substituting:

$$\underline{x} = \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} k & 0 \\ \frac{c}{d} & 1 \end{pmatrix} \begin{pmatrix} z \\ w \end{pmatrix} = \begin{pmatrix} kz \\ \frac{c}{d}z + w \end{pmatrix}$$

Similarly, it is possible to calculate  $\underline{z}$ :

$$\underline{z} = \begin{pmatrix} z \\ w \end{pmatrix} = \begin{pmatrix} \frac{1}{k} & 0 \\ -\frac{c}{d} \frac{1}{k} & 1 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} \frac{1}{k}x \\ -\frac{c}{d} \frac{1}{k}x + y \end{pmatrix}$$



## *Center manifold analysis in SISUV*

The ODE system from the original  $(I, V)$  to the  $\underline{z}$  system is given by  $\frac{d}{dt}\underline{z} = A\underline{z} + \underline{q}$  with the nonlinear part given by  $\underline{q} := zw \cdot \begin{pmatrix} -\frac{\beta}{M} \\ -\frac{\vartheta}{N} \end{pmatrix}$ . Now we can obtain the time derivative of the vector  $\underline{x}$  via:

$$\frac{d}{dt}\underline{x} = \Lambda\underline{x} + T^{-1}\underline{q}(\underline{x})$$

Obtaining explicitly:

$$\begin{aligned} \dot{x} &= -\frac{\beta}{M}xy + \frac{c}{dk}x^2 \\ \dot{y} &= d \cdot y + \left( \frac{c}{dM} \frac{\beta}{M} + \frac{\vartheta}{N} \right) \left( \frac{c}{dk}x^2 - \frac{1}{k}xy \right) \end{aligned}$$

To find the transformation  $y = h(x)$  along the center manifold, the functional  $\mathcal{N}(h(x))$  has to vanish:

$$\mathcal{N}(h(x)) = \frac{dh}{dx} \cdot f(x, h(x)) - (d \cdot h(x) + g(x, h(x))) = 0$$

## *Center manifold analysis in SISUV*

This equation can be solved via polynomial approximation of  $h(x)$  :

$$h(x) := a_2 \cdot x^2 + a_3 \cdot x^3 + a_4 \cdot x^4 + a_5 \cdot x^5 + \mathcal{O}(x^6)$$

The center manifold was calculated by a 3<sup>rd</sup> order polynomial:

$$a_2 = -\frac{c \cdot s}{d^2 \cdot k^2}$$

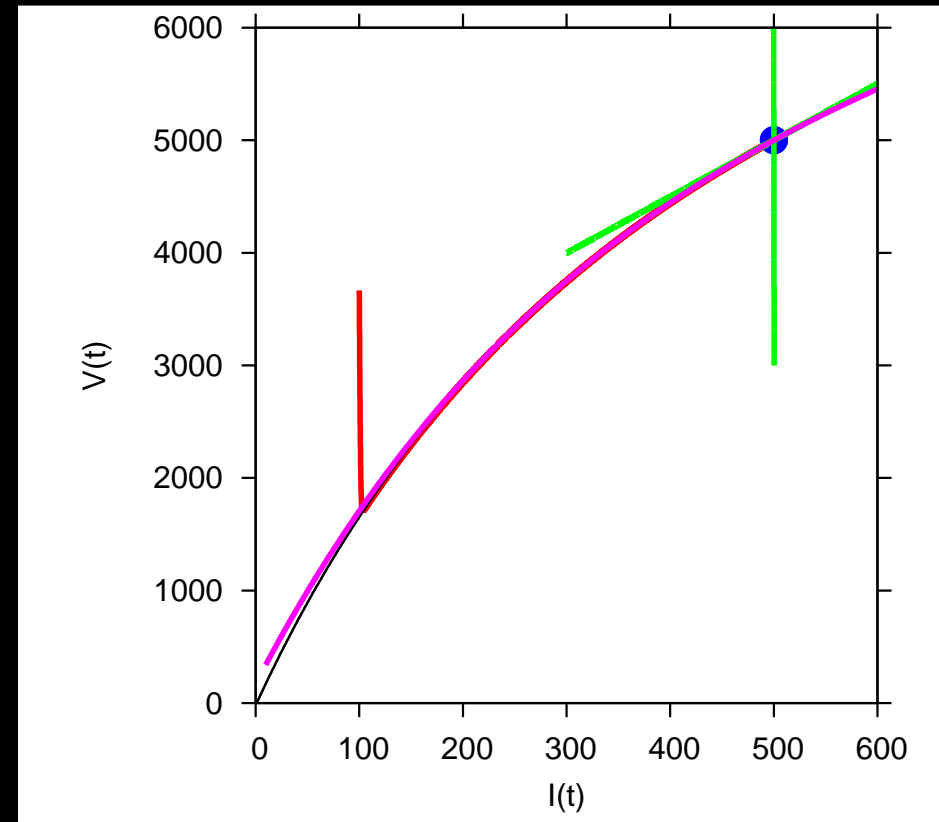
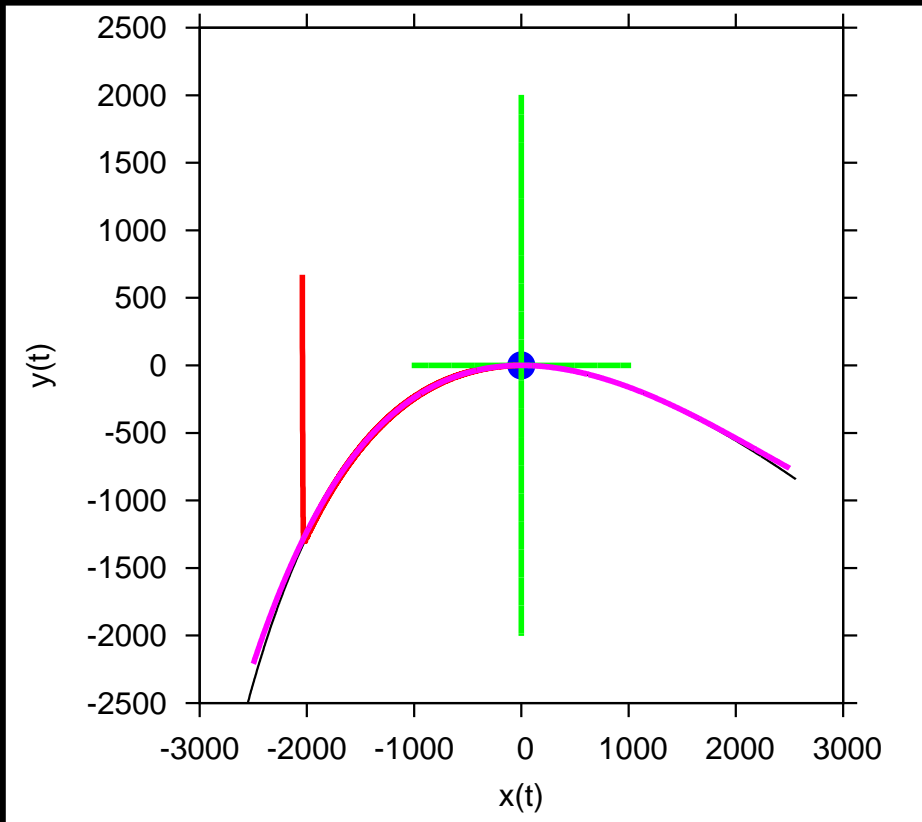
$$a_3 = \frac{1}{d} \left( \frac{2c}{d \cdot k} \frac{\beta}{M} + \frac{s}{k} \right) a_2$$

From the 3<sup>rd</sup> order polynomial, it is possible to use a general formula to easily get a polynomial of a higher order:

$$a_j = \frac{1}{d} \left( (j-1) \frac{\beta}{M} \frac{c}{k \cdot d} + \frac{s}{k} \right) a_{j-1} - \frac{\beta}{M \cdot d} \left( \sum_{\ell=2}^{j-2} \ell \cdot a_j \cdot a_{j-\ell} \right)$$

For  $j = 4, 5, \dots, \infty$ .

# *Center manifold analysis in SISUV*



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We started the analysis by the simplest SIS, because we can easily get the analytic solution for the model seasonal forced.

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Considering stable population size  $N = S + I$  we can simplify

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The seasonal forcing is given by  $\beta(t) = \beta_0(1 + \eta \cdot \cos(\omega t))$

$$\dot{I} = \frac{\beta(t)}{N}(N - I)I - \alpha I$$

## *Analytic seasonal forced SIS*

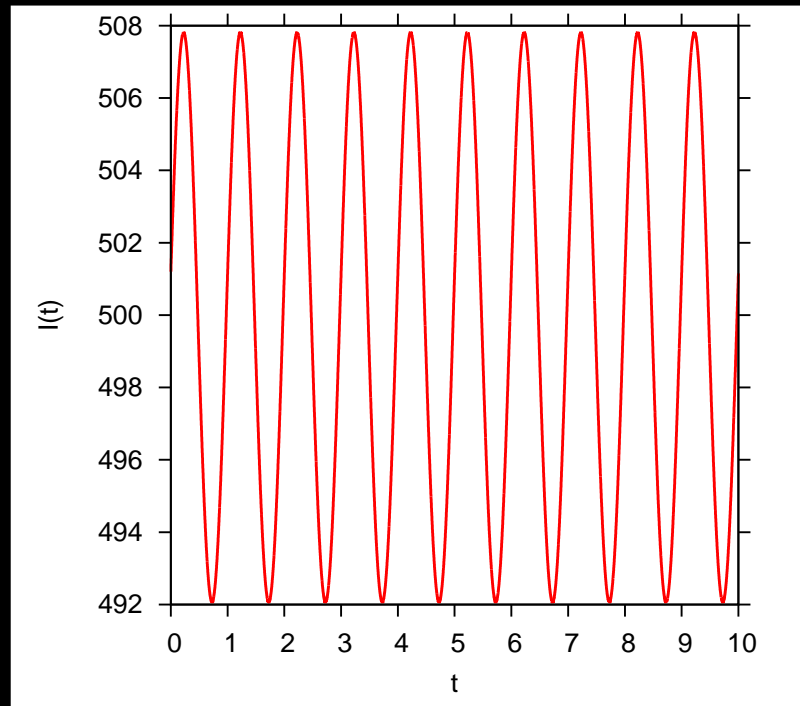
In the seasonal forcing we will consider the complex formulation, for now

$$\beta(t) = \beta_0 + \varepsilon\beta_1 e^{i\omega t}$$



# *Analytic seasonal forced SIS*

If we plot the SIS seasonal forced



The  $I(t)$  is defined by the stationary state plus some oscillations dependent on the amplitude  $I_1$ , *i.e.*

$$I(t) = I_0 + \varepsilon I_1 e^{i\omega t} + \mathcal{O}(\varepsilon^2)$$

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And applying the time derivative to  $I(t)$

$$\frac{dI}{dt} = \varepsilon I_1 i\omega e^{i\omega t}$$

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Substituting in the ODE

$$\begin{aligned}\frac{d}{dt}I &= \frac{\beta(t)}{N} (N - I) I - \alpha I \\ \varepsilon i\omega I_1 e^{i\omega t} &= \frac{1}{N} (\beta_0 + \varepsilon\beta_1 e^{i\omega t}) (N - (I_0 + \varepsilon I_1 e^{i\omega t})) (I_0 + \varepsilon I_1 e^{i\omega t}) - \alpha (I_0 + \varepsilon I_1 e^{i\omega t})\end{aligned}$$

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And separating the terms in respect to order of  $\varepsilon$ , we get

$$\begin{aligned}\varepsilon i\omega I_1 e^{i\omega t} &= \varepsilon e^{i\omega t} \left( -\alpha I_1 + \frac{1}{N} (-\beta_0 I_0 I_1 + I_0 \beta_1 N - I_0^2 \beta_1 + I_1 \beta_0 N - \beta_0 I_0 I_1) \right) \\ &\quad + \varepsilon^0 \left( \frac{\beta_0}{N} (N - I_0) I_0 - \alpha I_0 \right)\end{aligned}$$

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The values of order  $\varepsilon^0$  have conditions for stationarity, hence  $I_0 = I^*$

$$\varepsilon i\omega I_1 e^{i\omega t} = \varepsilon e^{i\omega t} \left( -\alpha I_1 + \frac{1}{N} (-\beta_0 I_0 I_1 + I_0 \beta_1 N - I_0^2 \beta_1 + I_1 \beta_0 N - \beta_0 I_0 I_1) \right)$$

## *Analytic seasonal forced SIS*

We get the complex amplitude for  $I_1$

$$I_1 = \frac{\frac{\beta_1}{N} (N - I_0) I_0}{i\omega + \frac{\beta_0}{N} I_0 + \alpha - \frac{\beta_0}{N} (N - I_0)}$$

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Setting  $a := \frac{\beta_0}{N} I_0 + \alpha - \frac{\beta_0}{N} (N - I_0)$  and  $c := \frac{\beta_1}{N} (N - I_0) I_0$ , we can simplify

$$I_1 = \frac{c}{a + i\omega}$$



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$$I_1 = \frac{c}{a + i\omega}$$

And multiplying numerator and denominator by its complex conjugate  $a - i\omega$

$$I_1 = \frac{ca}{(a^2 + \omega^2)} + i \left( \frac{-ca}{(a^2 + \omega^2)} \right)$$

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$$I_1 = \frac{ca}{(a^2 + \omega^2)} + i \left( \frac{-ca}{(a^2 + \omega^2)} \right) := \tilde{I}_1 + i\hat{I}_1$$

where the real part  $\tilde{I}_1 := \frac{ca}{(a^2 + \omega^2)}$  and the imaginary part  $\hat{I}_1 := \frac{-c\omega}{(a^2 + \omega^2)}$  are determined.

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Hence the complex response of  $I(t)$  is given by

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And applying the same calculations for  $e^{-i\omega t}$ , the second part of the real *cos* function for  $I_1$  using its complex conjugate  $\bar{I}_1 = \tilde{I}_1 - i\hat{I}_1$

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Combining the results for  $e^{i\omega t}$  and  $e^{-i\omega t}$  gives for the real forcing  $\beta(t) = \beta_0 + \varepsilon \frac{1}{2} \beta_1 (e^{i\omega t} + e^{-i\omega t})$  the real response of the infected

$$I(t) = I^* + \varepsilon \cdot A_I \cdot \cos(\omega(t + \varphi_I))$$

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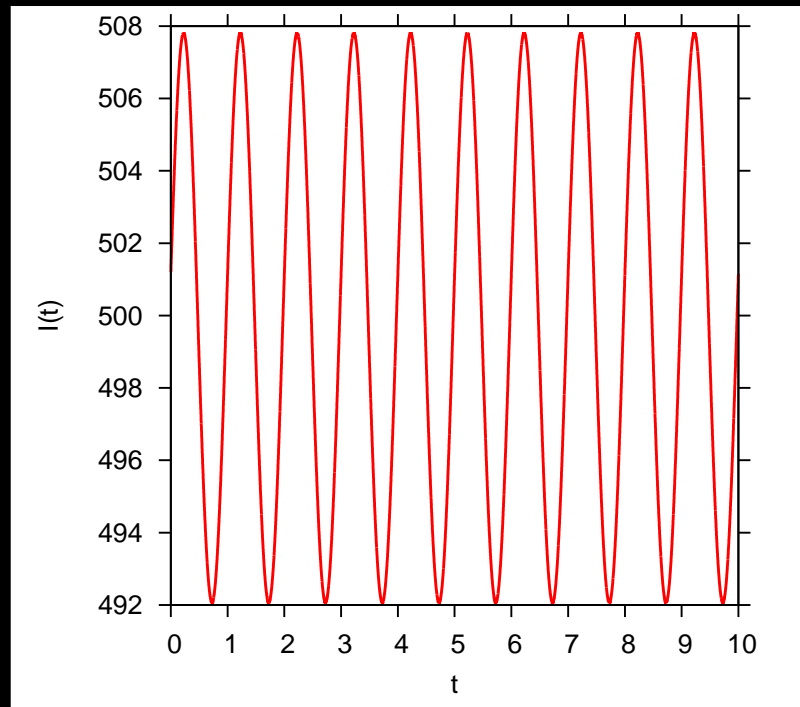
$$I(t) = I^* + \varepsilon \cdot A_I \cdot \cos(\omega(t + \varphi_I))$$

with real amplitude  $A_I$  and phase  $\varphi_I$  calculated from the complex amplitude

$$A_I = 2\sqrt{\tilde{I}_1^2 + \hat{I}_1^2}$$

$$\varphi_I = \frac{1}{\omega} \arctan\left(\frac{\hat{I}_1}{\tilde{I}_1}\right)$$

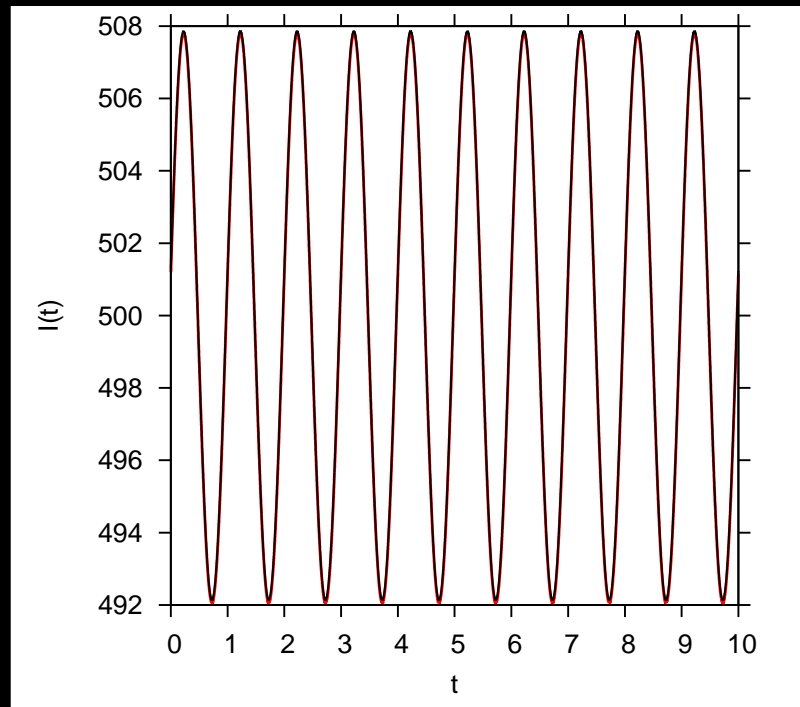
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Parameters:  $\alpha = \frac{1}{10}y^{-1}$ ,  $\beta_0 = 2\alpha$  and  $\eta = 0.1$



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The next step was to introduce the vector dynamic into the SIS system, getting the SISUV

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$$\dot{I} = \frac{\beta}{M}SV - \alpha I$$

$$\dot{U} = \psi - \nu U - \frac{\vartheta}{N}UI$$

$$\dot{V} = \frac{\vartheta}{N}UI - \nu V$$

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$$\dot{I} = \frac{\beta}{M}SV - \alpha I$$

$$\dot{U} = \psi - \nu U - \frac{\vartheta}{N}UI$$

$$\dot{V} = \frac{\vartheta}{N}UI - \nu V$$

Considering  $N = S(t) + I(t)$  and  $M = U(t) + V(t)$  we obtain:

$$\frac{d}{dt}I = \frac{\beta}{M}(N - I)V - \alpha I$$

$$\frac{d}{dt}V = \frac{\vartheta}{N}(M - V)I - \nu V$$

## *Analytic seasonally forced SISUV*

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$$\begin{aligned}\dot{S} &= \alpha I - \frac{\beta}{M}SV \\ \dot{I} &= \frac{\beta}{M}SV - \alpha I \\ \dot{U} &= \vartheta - \nu U - \frac{\vartheta}{N}UI \\ \dot{V} &= \frac{\vartheta}{N}UI - \nu V\end{aligned}$$

Considering  $N = S(t) + I(t)$  and  $M = U(t) + V(t)$  we obtain:

$$\begin{aligned}\frac{d}{dt}I &= \frac{\beta}{M}(N - I)V - \alpha I \\ \frac{d}{dt}V &= \frac{\vartheta}{N}(M(t) - V)I - \nu V\end{aligned}$$

With the seasonal forcing given by  $M(t) = M_0(1 + \eta \cdot \cos(\omega t))$ .

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The real part of the seasonal forcing is

$$M(t) = M_0 + \varepsilon M_1 e^{i\omega t}$$

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$$I(t) = I_0 + \varepsilon I_1 e^{i\omega t} + \mathcal{O}(\varepsilon^2)$$

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and the time derivatives

$$\frac{dI}{dt} = \varepsilon I_1 i\omega e^{i\omega t}$$

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## *Analytic seasonally forced SISUV*

Substituting in the ODE for  $I$

$$\frac{d}{dt}I = \frac{\beta}{M}(N - I)V - \alpha I$$

$$\varepsilon I_1 i \omega e^{i\omega t} I = \frac{\beta}{M} (N - (I_0 + \varepsilon I_1 e^{i\omega t})) (V_0 + \varepsilon V_1 e^{i\omega t}) - \alpha (I_0 + \varepsilon I_1 e^{i\omega t})$$



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And reorganizing the terms of different orders of  $\varepsilon$

$$\varepsilon I_1 i \omega e^{i \omega t} I = \frac{\beta}{M} (N - I_0) V_0 - \alpha I_0 + \varepsilon e^{i \omega t} \left[ \frac{\beta}{M} (N V_1 - I_0 V_1 - V_0 I_1) - \alpha I_1 \right]$$

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As  $I_0 = I^*$  and  $V_0 = V^*$  we can say that  $\frac{\beta}{M} (N - I_0) V_0 - \alpha I_0 = 0$ , so:

$$\varepsilon I_1 i \omega e^{i\omega t} = \varepsilon e^{i\omega t} \left[ \frac{\beta}{M} (N V_1 - I_0 V_1 - V_0 I_1) - \alpha I_1 \right]$$

## *Analytic seasonally forced SISUV*

And finally we get the complex amplitude of  $I_1$  dependent on  $V_1$

$$I_1 = \frac{\frac{\beta}{M} (N - I_0)}{\frac{\beta}{M} V_0 + \alpha + i\omega} V_1$$

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$$I_1 = \frac{\frac{\beta}{M}(N - I_0)}{\frac{\beta}{M}V_0 + \alpha + i\omega} V_1$$

Setting  $\frac{\beta}{M}(N - I_0) := c$  and  $\frac{\beta}{M}V_0 + \alpha := d$ , we get

$$I_1 := \frac{c}{d + i\omega} V_1$$

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And multiplying numerator and denominator by the complex conjugate  $\frac{c}{d - i\omega}$ , we obtain

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$$I_1 = \left( \frac{cd}{d^2 + \omega^2} + i \frac{-c\omega}{d^2 + \omega^2} \right) V_1 =: (a + ib) V_1$$

with  $a := \frac{cd}{d^2 + \omega^2}$  and  $b := \frac{-c\omega}{d^2 + \omega^2}$ .

## *Analytic seasonally forced SISUV*

Now we apply the same calculations to find the analytic solution of the amplitude for  $V_1$ . However, in this case we use the seasonal forcing in  $M(t)$

$$\frac{d}{dt}V = \frac{\vartheta}{N}(M(t) - V)I - \nu V$$

$$\begin{aligned} \varepsilon V_1 i\omega e^{i\omega t} &= \frac{\vartheta}{N} \left( (M_0 + \varepsilon M_1 e^{i\omega t}) - (V_0 + \varepsilon V_1 e^{i\omega t}) \right) (V_0 + \varepsilon (a + ib) V_1 e^{i\omega t}) - \\ &\quad - \nu (V_0 + \varepsilon V_1 e^{i\omega t}) \end{aligned}$$

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Rearranging in orders of  $\varepsilon$ , we get

$$\begin{aligned} \varepsilon V_1 i\omega e^{i\omega t} &= \frac{\vartheta}{N} (M_0 - V_0) I_0 - \nu V_0 + \\ &\quad + \varepsilon e^{i\omega t} \left[ \frac{\vartheta}{N} M_1 I_0 + \left( \frac{\vartheta}{N} (M_0(a + ib) - I_0 - V_0(a + ib)) - \nu \right) V_1 \right] \end{aligned}$$



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Now we apply the same calculations to find the analytic solution of the amplitude for  $V_1$ . However, in this case we use the seasonal forcing in  $M(t)$

$$\frac{d}{dt}V = \frac{\vartheta}{N}(M(t) - V)I - \nu V$$

$$\begin{aligned} \varepsilon V_1 i \omega e^{i\omega t} &= \frac{\vartheta}{N} \left( (M_0 + \varepsilon M_1 e^{i\omega t}) - (V_0 + \varepsilon V_1 e^{i\omega t}) \right) (V_0 + \varepsilon (a + ib) V_1 e^{i\omega t}) - \\ &\quad - \nu (V_0 + \varepsilon V_1 e^{i\omega t}) \end{aligned}$$

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Once again, we can forget about the terms of the order  $\varepsilon^0$

$$V_1 i \omega = \frac{\vartheta}{N} M_1 I_0 + \left( \frac{\vartheta}{N} (M_0 (a + ib) - I_0 - V_0 (a + ib)) - \nu \right) V_1$$

## *Analytic seasonally forced SISUV*

Obtaining specifically

$$V_1 = \frac{\frac{\vartheta}{N} M_1 I_0}{\frac{\vartheta}{N} (I_0 + a(V_0 - M_0)) + \nu + i \left( \frac{\vartheta}{N} b(V_0 - M_0) + \omega \right)}$$

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Setting  $u := \frac{\vartheta}{N} (I_0 + a(V_0 - M_0)) + \nu$ ,  $v := \frac{\vartheta}{N} b(V_0 - M_0) + \omega$  and  $w := \frac{\vartheta}{N} M_1 I_0$ .

$$V_1 =: \frac{w}{u + iv}$$

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$$V_1 =: \frac{w}{u + iv}$$

Multiplying both terms by the complex conjugate, we get

$$V_1 = \frac{wu}{u^2 + v^2} + i \frac{-wv}{u^2 + v^2} =: \tilde{V}_1 + i\hat{V}_1$$

being  $\tilde{V}_1 := \frac{wu}{u^2 + v^2}$  and  $\hat{V}_1 := \frac{-wv}{u^2 + v^2}$ , obtaining the complex amplitude for  $V_1$ .

## *Analytic seasonally forced SISUV*

And substituting in the analytic expression for complex amplitude for  $I_1$ , we obtain

$$I_1 = (a + ib) \cdot (\tilde{V}_1 + i\hat{V}_1)$$

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And setting  $\tilde{I}_1 := a\tilde{V}_1 - b\hat{V}_1$  and  $\hat{I}_1 := a\hat{V}_1 + b\tilde{V}_1$ , we can simplify

$$I_1 =: \tilde{I}_1 + i\hat{I}_1$$

## *Analytic seasonally forced SISUV*

Obtaining the complex response for both  $I(t)$  and  $V(t)$  with the real and complex parts of each one and doing a similar analysis for the second part of the real  $\cos$  function, as was done for SIS, we can obtain the real response for the system:



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$$I(t) = I^* + \varepsilon \cdot A_I \cdot \cos(\omega(t + \varphi_I))$$

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$$I(t) = I^* + \varepsilon \cdot A_I \cdot \cos(\omega(t + \varphi_I))$$

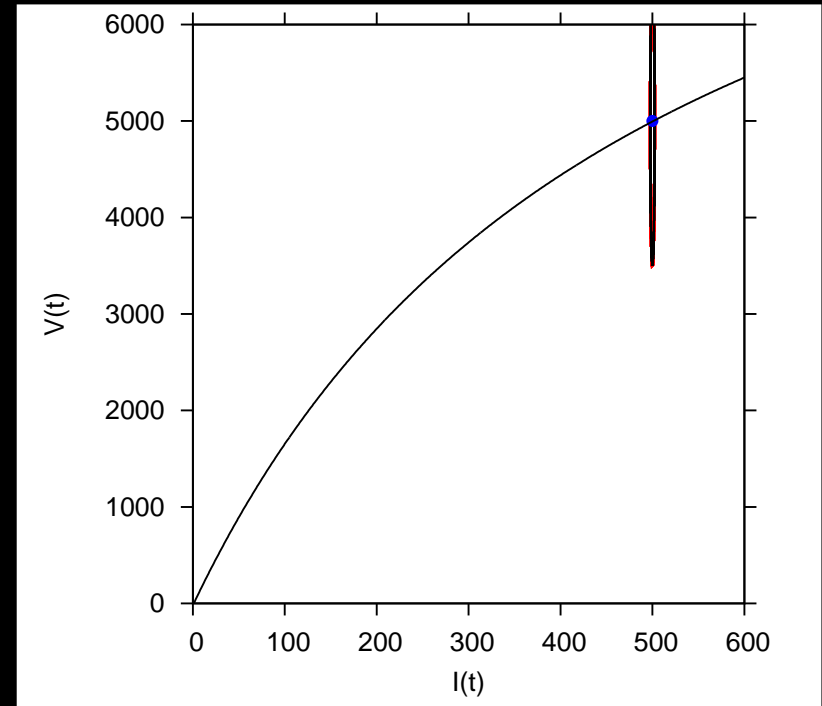
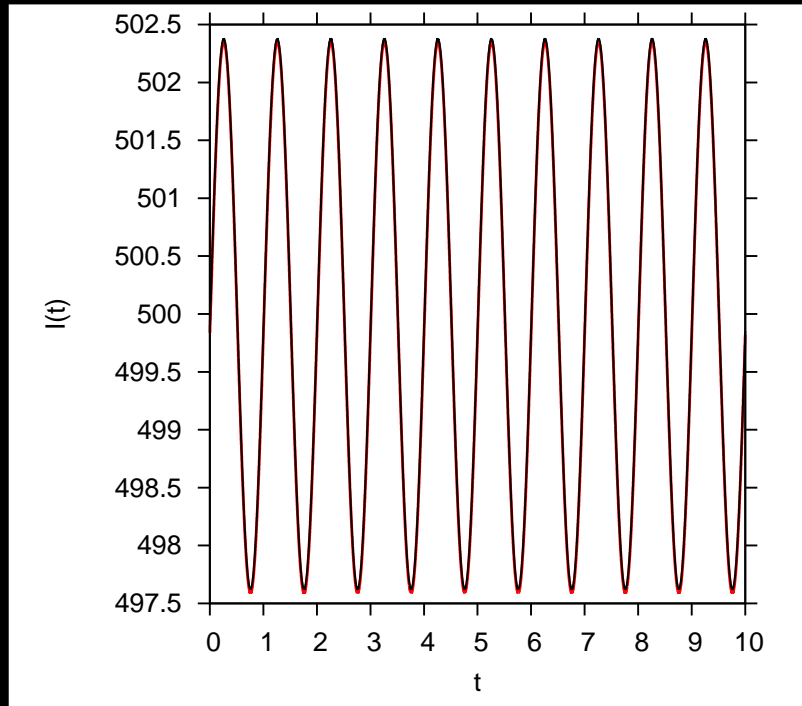
$$V(t) = V^* + \varepsilon \cdot A_V \cdot \cos(\omega(t + \varphi_V))$$

with the amplitude and phase for both variables given by

$$A_I = 2\sqrt{\tilde{I}_1^2 + \hat{I}_1^2} \quad \varphi_I = \frac{1}{\omega} \arctan\left(\frac{\hat{I}_1}{\tilde{I}_1}\right)$$

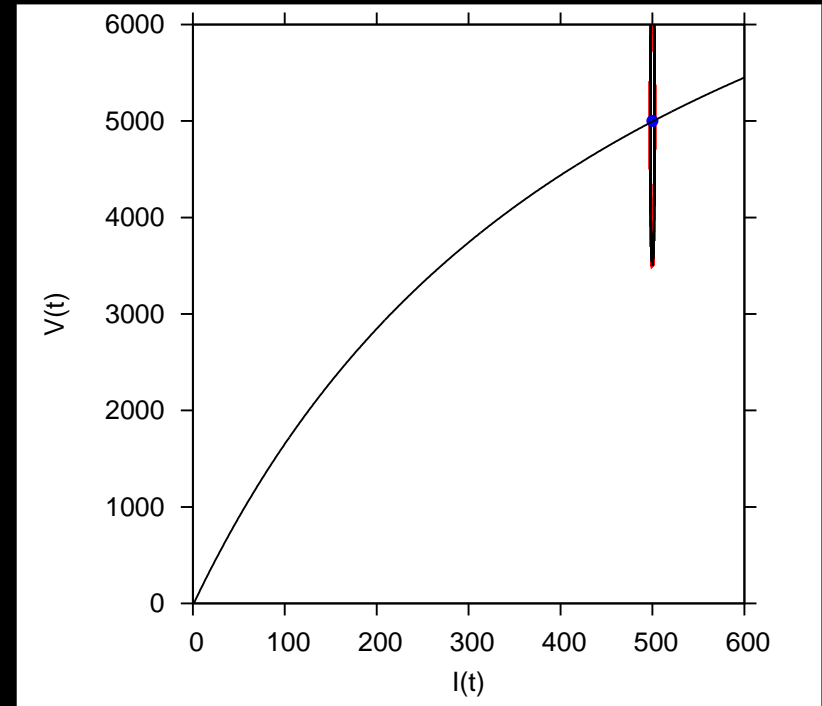
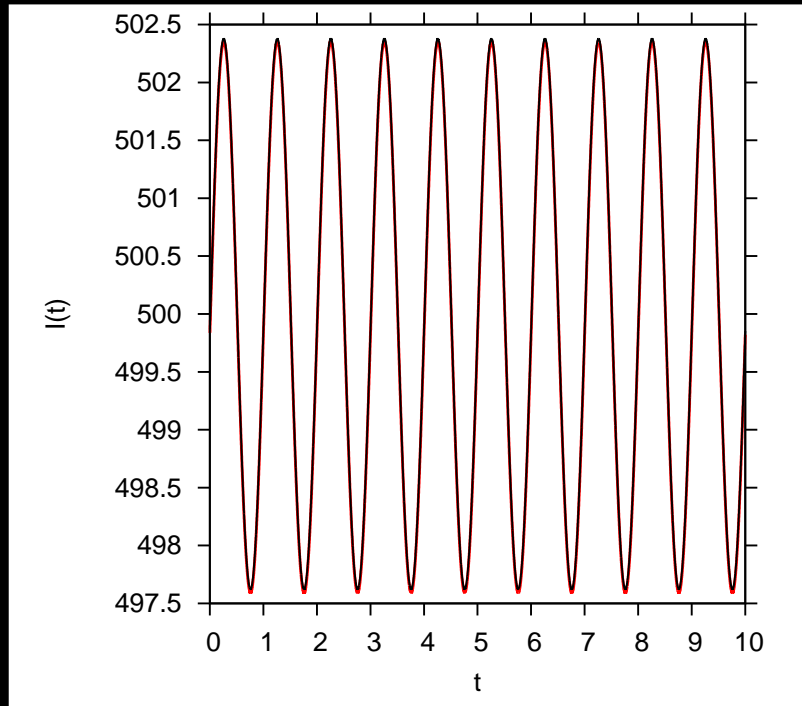
$$A_V = 2\sqrt{\tilde{V}_1^2 + \hat{V}_1^2} \quad \varphi_V = \frac{1}{\omega} \arctan\left(\frac{\hat{V}_1}{\tilde{V}_1}\right)$$

# *Analytic seasonally forced SISUV*



Parameters:  $\alpha = \frac{1}{10}y^{-1}$  ,  $\beta = 2\alpha$  ,  $\nu = \frac{365}{10}d^{-1}$  ,  $\vartheta = 2\nu$  and  $\eta = 0.3$  .

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The amplitude in  $V(t)$  caused by the seasonality is not reflected in  $I(t)$  dynamics.

So, for modelling proposes, the vector dynamics is not important for the system.

## *The Full SIRUV and comparison with SIR model*

After have analysed the simplest models we made similar calculation for more complicated SIR and SIRUV model, comparing the expressions from each other.

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Considering a closed population for humans  $N(t) = S(t) + I(t) + R(t)$  we can simplify the SIR model into a two dimensional system:

$$\frac{d}{dt}I = \frac{\beta(t)}{N} \cdot (N - I - R) \cdot I - (\mu + \gamma) \cdot I$$

$$\frac{d}{dt}R = \gamma \cdot I - \mu \cdot R$$

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And also for vectors  $M(t) = U(t) + V(t)$  for SIRUV model

$$\frac{d}{dt}I = \frac{\beta}{M_0} \cdot (N - I - R) \cdot V - (\gamma + \mu) \cdot I$$

$$\frac{d}{dt}R = \gamma \cdot I - \mu \cdot R$$

$$\frac{d}{dt}V = \frac{\vartheta}{N} \cdot (M(t) - V) \cdot I - \nu \cdot V \quad .$$

## *The Full SIRUV and comparison with SIR model*

The endemic stationary states of the SIR are given by

$$I^* = \frac{\mu}{(\gamma + \mu)} \cdot \left(1 - \frac{\gamma + \mu}{\beta}\right) \cdot N$$
$$R^*(I^*) = \frac{\gamma}{\mu} I^*$$



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Whereas the stationary states for SIRUV system are given by

$$I^* = \frac{\mu}{(\gamma + \mu) + \frac{\gamma + \mu}{\beta} \cdot \mu} \cdot \left(1 - \frac{\gamma + \mu}{\frac{\vartheta}{\nu} \beta}\right) \cdot N$$
$$R^*(I^*) = \frac{\gamma}{\mu} I^*$$
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$$I^* = \frac{\mu}{(\gamma + \mu) + \frac{\gamma + \mu}{\beta} \cdot \mu} \cdot \left(1 - \frac{\gamma + \mu}{\frac{\vartheta}{\nu}\beta}\right) \cdot N$$
$$R^*(I^*) = \frac{\gamma}{\mu} I^*$$
$$V^*(I^*) = \frac{\frac{\vartheta}{N} I^*}{\nu + \frac{\vartheta}{N} I^*} \cdot M \quad .$$

Essentially we can obtain the same stationary state in both models, by replacing the  $\beta$  by  $\frac{\vartheta}{\nu}\beta$  in SIR.

## *The Full SIRUV and comparison with SIR model*

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For small seasonal forcing  $\eta$ , hence small  $\varepsilon$  we expect also small oscillations of the state variables, hence

$$I(t) = I_0 + \varepsilon I_1 e^{i\omega t} + \mathcal{O}(\varepsilon^2)$$

$$R(t) = R_0 + \varepsilon R_1 e^{i\omega t} + \mathcal{O}(\varepsilon^2)$$

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and for the time derivatives

$$\frac{dI}{dt} = \varepsilon I_1 i\omega e^{i\omega t}$$

$$\frac{dR}{dt} = \varepsilon R_1 i\omega e^{i\omega t}$$

$$\frac{dV}{dt} = \varepsilon V_1 i\omega e^{i\omega t} .$$

## *The Full SIRUV and comparison with SIR model*

The  $R$  is the same for both models, so

$$\frac{d}{dt}R = \gamma I - \mu R$$

$$\varepsilon R_1 i \omega e^{i\omega t} = \gamma (I_0 + \varepsilon I_1 e^{i\omega t}) - \mu (R_0 + \varepsilon R_1 e^{i\omega t})$$

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To zeroth order  $\varepsilon^0$  we obtain again the condition for stationarity

$$R_1 = \frac{\gamma}{\mu + i\omega} I_1$$



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Finally, multiplying numerator and denominator both by its complex conjugate we obtain

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with  $a := \frac{\gamma\mu}{\mu^2 + \omega^2}$   $b := \frac{-\gamma\omega}{\mu^2 + \omega^2}$ .

## *The Full SIRUV and comparison with SIR model*

Now we are going to analyse the I for both models.

$$\frac{dI}{dt} = \frac{\beta(t)}{N} \cdot (N - I - R) \cdot I - (\mu + \gamma) \cdot I \quad \frac{dI}{dt} = \frac{\beta}{M} \cdot (N - I - R) \cdot V - (\gamma + \mu) \cdot I$$

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$$\frac{d}{dt}I = \frac{\beta(t)}{N} \cdot (N - I - R) \cdot I - (\mu + \gamma) \cdot I \qquad \frac{d}{dt}I = \frac{\beta}{M} \cdot (N - I - R) \cdot V - (\gamma + \mu) \cdot I$$

$$I_1 = \frac{\frac{\beta_1}{N}(N - I_0 - R_0)}{\frac{\beta_0}{N}(1+a)I_0 - \frac{\beta_0}{N}(N - I_0 - R_0) + (\gamma + \mu) + i\left(\omega + \frac{\beta_0}{N}bI_0\right)} I_0 \qquad I_1 = \frac{\frac{\beta}{M_0}(N - I_0 - R_0)}{\frac{\beta}{M_0}(1+a)V_0 + (\gamma + \mu) + i\left(\omega + \frac{\beta}{M_0}bV_0\right)} V_1$$

# *The Full SIRUV and comparison with SIR model*

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Putting both in a  $I_1 := \frac{f}{c+id}$  form, we have the specific values of each abbreviation.

$$c_{IR} = \frac{\beta_0}{N} (1 + a) I_0 - \frac{\beta_0}{N} (N - I_0 - R_0) + (\gamma + \mu) \quad c_{IRV} = \frac{\beta}{M_0} (1 + a) V_0 + (\gamma + \mu)$$

$$d_{IR} = \omega + \frac{\beta_0}{N} b I_0$$

$$d_{IRV} = \omega + \frac{\beta}{M_0} b V_0$$

$$f_{IR} = \frac{\beta_1}{N} (N - I_0 - R_0) I_0$$

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$$f_{IRV} = \frac{\beta}{M_0} (N - I_0 - R_0)$$

Essentially, the only difference between the models is given by  $-\frac{\beta_0}{N} (N - I_0 - R_0)$  in the SIR.

## *The Full SIRUV and comparison with SIR model*

Multiplying both terms of  $I_1$  by the complex conjugate  $c-id$  we obtain the amplitude of  $I_1$  for models.

$$I_1 := \frac{f_{IR}}{c_{IR} + id_{IR}}$$

$$I_1 := \frac{f_{IRV}}{c_{IRV} + id_{IRV}} V_1$$

$$I_1 = \left( \left( \frac{f_{IR} c_{IR}}{c_{IR}^2 + d_{IR}^2} \right) + i \left( \frac{-f_{IR} d_{IR}}{c_{IR}^2 + d_{IR}^2} \right) \right)$$

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$$I_1 := \tilde{I}_1 + i\hat{I}_1$$

$$I_1 := (x + iy)V_1$$

Being  $\tilde{I}_1 \triangleq x := \frac{fc}{c^2+d^2}$  and  $\hat{I}_1 \triangleq y := \frac{-fd}{c^2+d^2}$ .



## *The Full SIRUV and comparison with SIR model*

Finally, from the ODE for the infected mosquitos

$$\frac{d}{dt}V = \frac{\vartheta}{N}(M(t) - V)I - \nu V$$

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$$V_1 := \frac{k}{g + ih}$$

with the coefficients  $g := \frac{\vartheta}{N}(I_0 - (M_0 - V_0)x) + \nu$ ,  $h := \omega - \frac{\vartheta}{N}(M_0 - V_0)y$  and  $k := \frac{\vartheta}{N}M_1I_0$ .

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where the real part is  $\tilde{V}_1 := \frac{kg}{g^2+h^2}$  and the imaginary is  $\hat{V}_1 := \frac{-kh}{g^2+h^2}$ .

## *The Full SIRUV and comparison with SIR model*

We can now substitute the complex amplitude of  $V_1$  in the amplitude of  $I_1$

$$I_1 = (x + iy) \cdot (\tilde{V}_1 + i\hat{V}_1)$$

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## *The Full SIRUV and comparison with SIR model*

So we have already obtained the first part of the real *cos* function for the three variables

$$I(t) = I^* + \varepsilon \cdot (\tilde{I}_1 + i\hat{I}_1)e^{i\omega t}$$

$$R(t) = R^* + \varepsilon \cdot (\tilde{R}_1 + i\hat{R}_1)e^{i\omega t}$$

$$V(t) = V^* + \varepsilon \cdot (\tilde{V}_1 + i\hat{V}_1)e^{i\omega t}$$

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And for the second part of the real *cos* function

$$I(t) = I^* + \varepsilon \cdot (\tilde{I}_1 - i\hat{I}_1)e^{-i\omega t}$$

$$R(t) = R^* + \varepsilon \cdot (\tilde{R}_1 - i\hat{R}_1)e^{-i\omega t}$$

$$V(t) = V^* + \varepsilon \cdot (\tilde{V}_1 - i\hat{V}_1)e^{-i\omega t}$$

## *The Full SIRUV and comparison with SIR model*

Combining the results for  $e^{i\omega t}$  and  $e^{-i\omega t}$  we obtain the real response of each variable.

$$I(t) = I^* + \varepsilon \cdot A_I \cdot \cos(\omega(t + \varphi_I))$$

$$R(t) = R^* + \varepsilon \cdot A_R \cdot \cos(\omega(t + \varphi_R))$$

$$V(t) = V^* + \varepsilon \cdot A_V \cdot \cos(\omega(t + \varphi_V))$$

with the real amplitude  $A$  and real phase  $\varphi$  calculated from the complex amplitude, via

$$A_I = 2\sqrt{\tilde{I}_1^2 + \hat{I}_1^2} \quad \varphi_I = \frac{1}{\omega} \arctan\left(\frac{\hat{I}_1}{\tilde{I}_1}\right)$$

$$A_R = 2\sqrt{\tilde{R}_1^2 + \hat{R}_1^2} \quad \varphi_R = \frac{1}{\omega} \arctan\left(\frac{\hat{R}_1}{\tilde{R}_1}\right)$$

$$A_V = 2\sqrt{\tilde{V}_1^2 + \hat{V}_1^2} \quad \varphi_V = \frac{1}{\omega} \arctan\left(\frac{\hat{V}_1}{\tilde{V}_1}\right)$$

## *The Full SIRUV and comparison with SIR model*

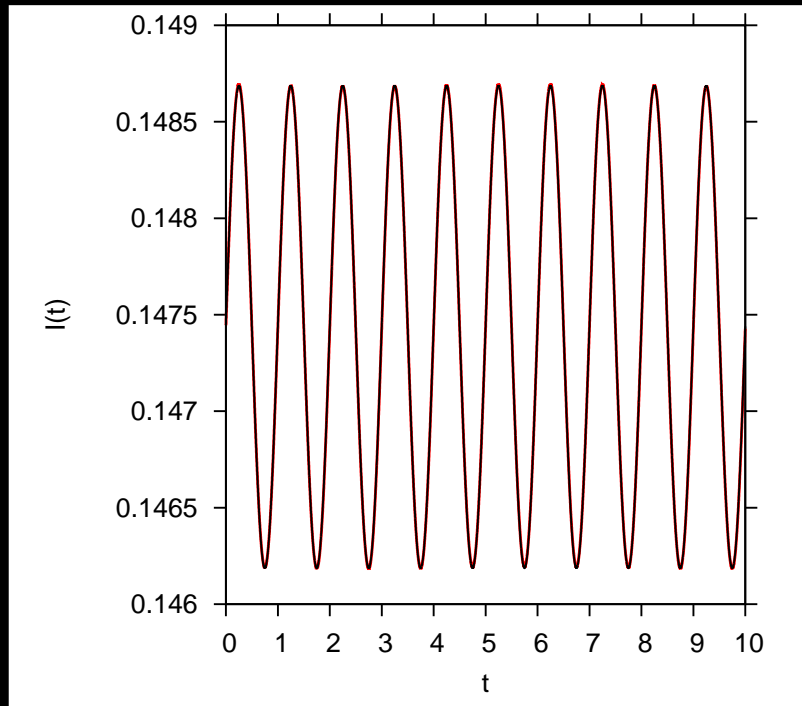
Comparing the amplitude and the phase numerically we have:

<i>SIR</i>	<i>SIRUV</i>
$\varepsilon A_I = 0.001955$	$\varepsilon A_I = 0.000774$
$\varphi_I = -0.248730$	$\varphi_I = -0.260094$

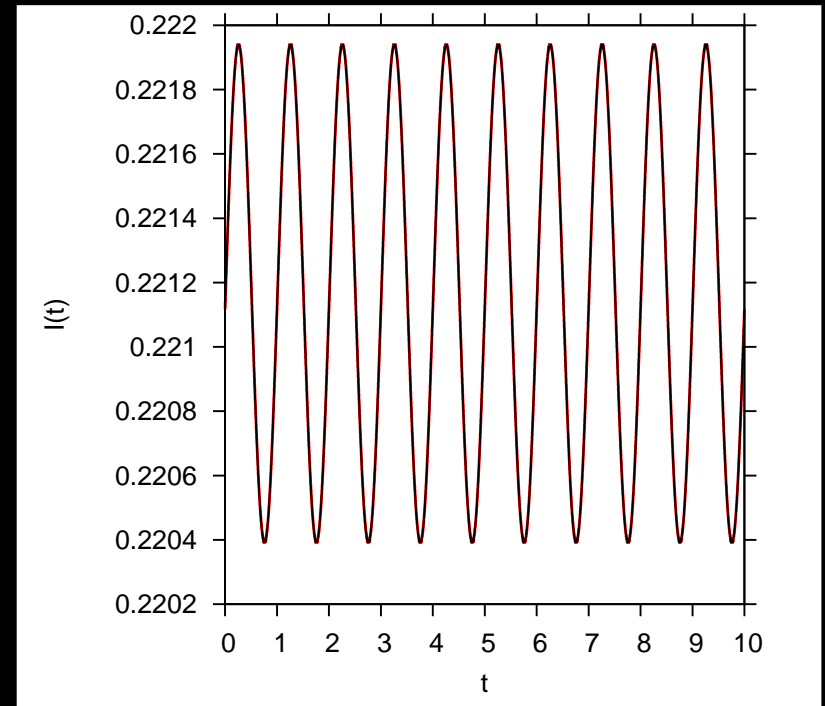
Essentially, the only difference is in the amplitude for both models.

# *The Full SIRUV and comparison with SIR model*

**SIR**



**SIRUV**

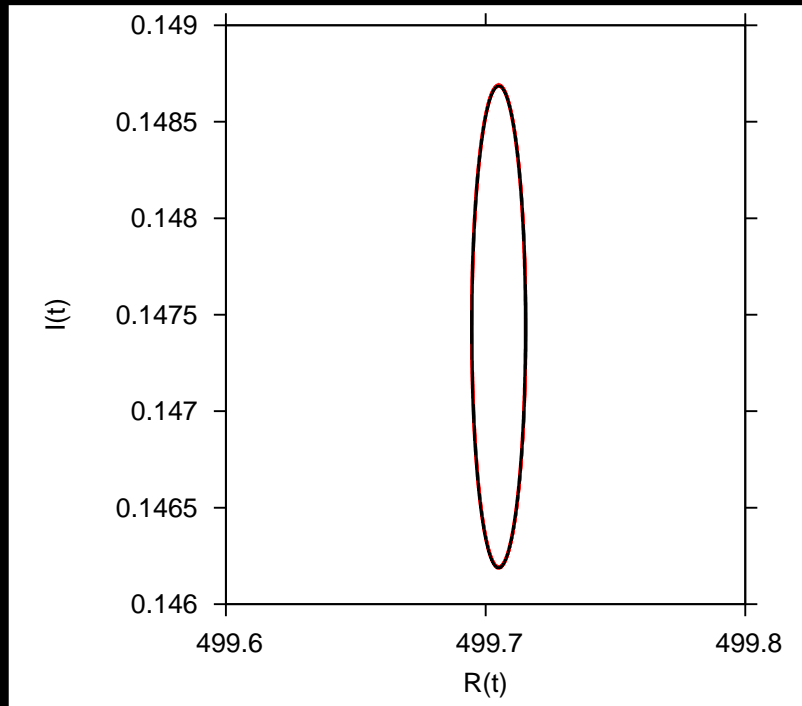


Parameters:  $\gamma = \frac{365}{7}d^{-1}$  ,  $\beta = 2\gamma$  ,  $\mu = \frac{1}{65}y^{-1}$  ,  $\nu = \frac{365}{10}d^{-1}$  ,  $\vartheta = 2\nu$  and  $\eta = 0.001$

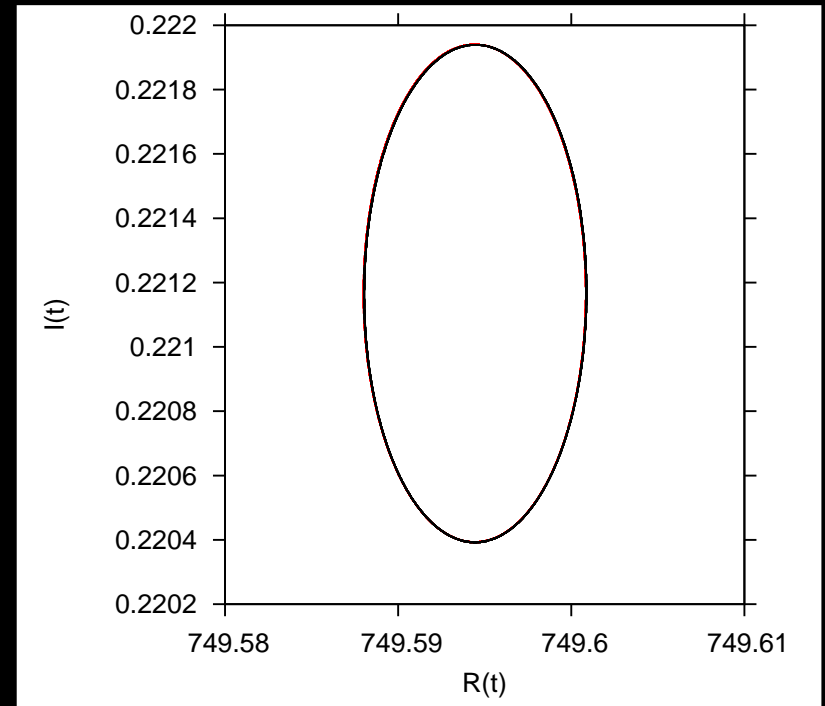


# *The Full SIRUV and comparison with SIR model*

**SIR**

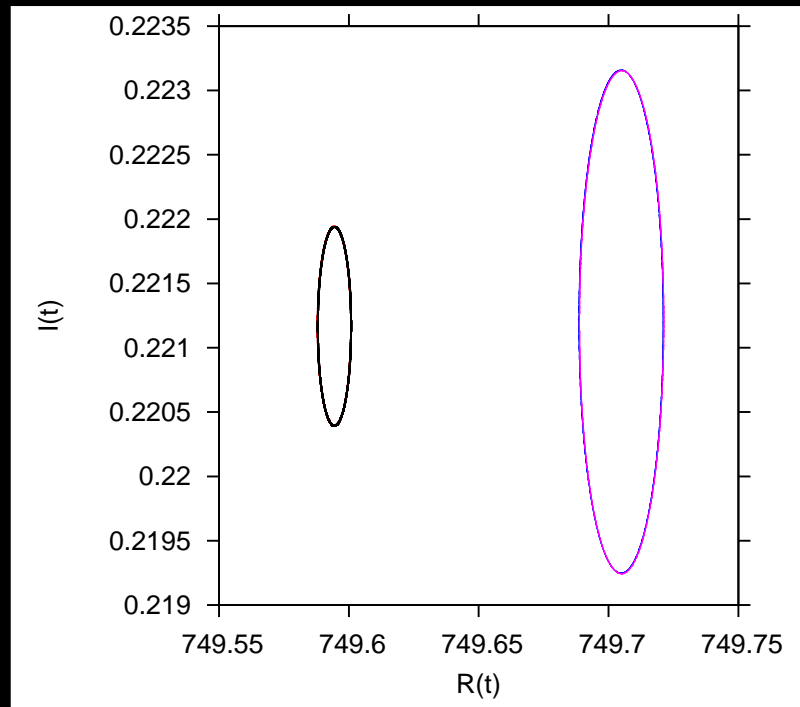


**SIRUV**



# *The Full SIRUV and comparison with SIR model*

We can join both models in the same plot and replace  $\beta$  of SIR by  $\frac{\vartheta}{\nu}\beta$  in order to have the stationary states approximated.

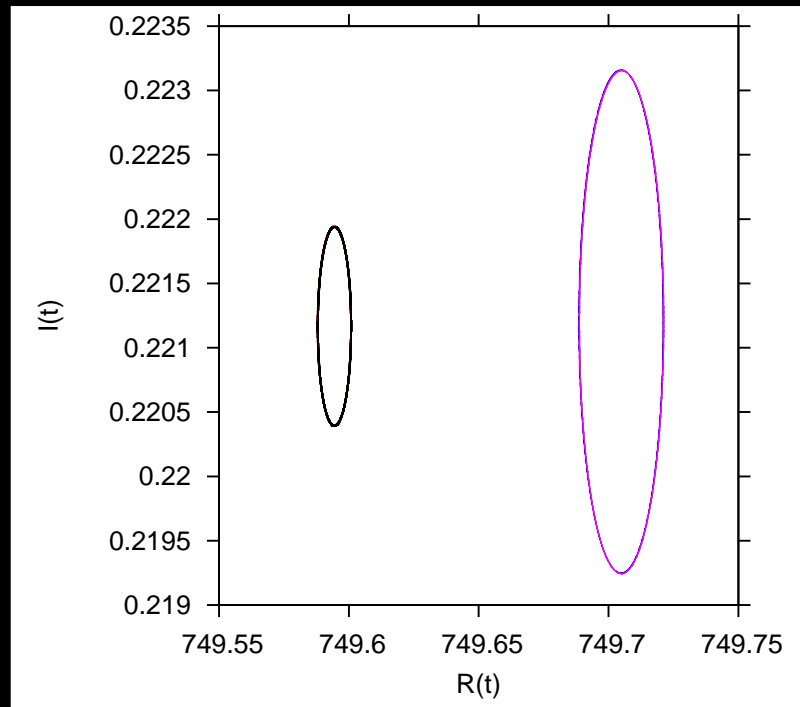


SIRUV

SIR

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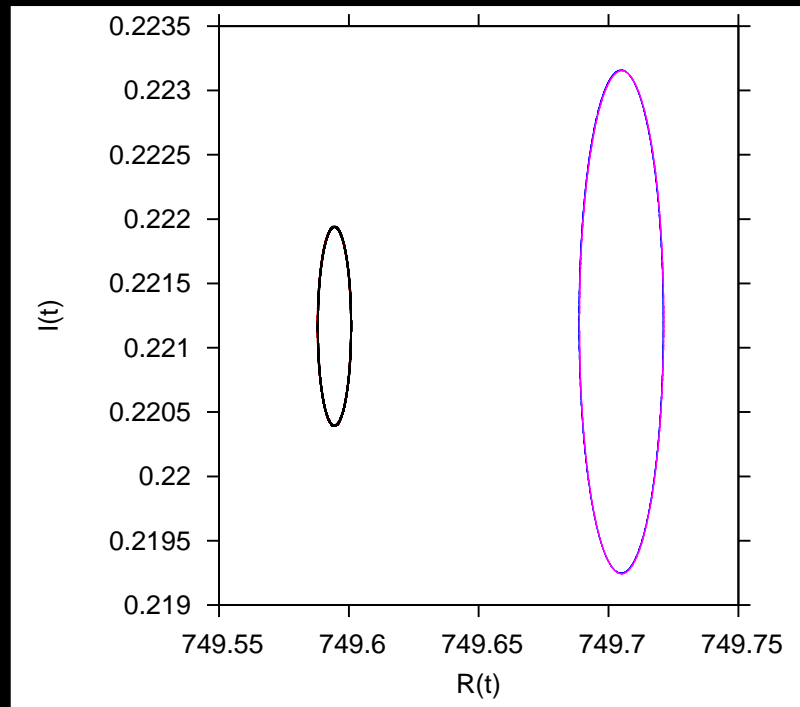
SIRUV

SIR

Comparing the two models we can say that are not such different, the differences in the amplitudes of  $I$  are in order of 0.001, so it is basically the same for both models.

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SIRUV

SIR

Comparing the two models we can say that are not such different, the differences in the amplitudes of  $I$  are in order of 0.001, so it is basically the same for both models.

So, once more, we can say that mosquitos do not add any information to models.

*Thank you for your attention!*

